

Estimating Lagrangian transport blending drifters with HF radar data and models: Results from the TOSCA experiment in the Ligurian Current (North Western Mediterranean Sea)

Maristella Berta^{a,*}, Lucio Bellomo^b, Marcello G. Magaldi^{a,c}, Annalisa Griffa^{a,d}, Anne Molcard^b, Julien Marmain^b, Mireno Borghini^a, Vincent Taillandier^e

^a*CNR-ISMAR, La Spezia, Italy*

^b*MIO, Université du Sud Toulon-Var, Aix-Marseille Université, CNRS/INSU, IRD, UM 110, La Garde, France*

^c*Johns Hopkins University, Baltimore, USA*

^d*MPO-RSMAS, University of Miami, Miami, USA*

^e*Laboratoire d'Océanographie de Villefranche, CNRS/UPMC, Villefranche sur Mer, France*

Abstract

Lagrangian transport estimates are investigated using results from HF radar, model and drifter data during a dedicated experiment in the Ligurian Current in the Toulon area (North Western Mediterranean Sea). Uncertainty estimates on particle position, $D(t)$, are computed and compared to absolute dispersion, $D_0(t)$, that provides an indication of the uncertainty in case of zero prior knowledge. In agreement with previous studies, radar results show that $D(t) \sim 1/2D_0(t)$ (i.e. ~ 6 km after 24 h). Model results are less reliable, as it can be expected in highly nonlinear coastal flows without local data assimilation. The central result of this paper is that when drifters are promptly deployed in an area of interest, their data can be used to significantly improve transport estimates using the Lagrangian blending algorithm LAVA with velocity fields from models or radar. Uncertainty can be reduced to $\sim 1/6D_0(t)$, (i.e. ~ 2 km after 24 h) for both radar and model, implying a much reduced search range in

*CNR-ISMAR, Forte Santa Teresa, Pozzuolo di Lerici, 19032 (SP). Italy. Tel: (+39)0187-1788916

Email address: maristella.bera@sp.ismar.cnr.it (Maristella Berta)

case of operational applications. The method is also found to have some forecasting skills with uncertainty $\sim 1/2D_0(t)$ during the first ~ 6 hours. Sensitivity tests provide indications on relevant time and space scales of predictability and provide suggestions for appropriate drifter sampling strategies.

Keywords: LAgrangian Variational Analysis (LAVA), HF radar, drifters, Lagrangian assimilation, North Western Mediterranean Sea, Ligurian-Provençal Current.

1. Introduction

Estimates of Lagrangian transport, i.e. transport of water following quantities carried by the marine currents, play a major role in many oceanographic applications. These include search and rescue (SAR) problems, management and mitigation of oil spills and other pollutant dispersals, as well as evaluation of sediment or larval transport for coastal or fishery management. While forecasts of ocean currents are desirable for all these applications, nowcasts are also very useful especially in case of SAR or of pollutant dispersals that cannot be easily observed at all times.

Evaluating Lagrangian transport is very challenging since the motion of particles advected by ocean currents is typically chaotic, i.e. highly dependent on initial conditions and on the details of the flow [Griffa et al., 2004]. In the last decade, results from instruments such as HF coastal radars and drifters have been used in many Lagrangian transport studies of the ocean surface [Molcard et al., 2009; Shadden et al., 2009; Ullman et al., 2006]. These two types of instruments provide different and in many respects complementary information. HF radars provide maps of surface velocity with ranges up to 100 km, horizontal resolution of the order 1.5-3 km, and temporal resolution of the order of 0.25-1 h [Gurgel et al., 1999b; Harlan et al., 2010]. Drifters, on the other hand, are water following instruments, providing direct and localized information on transport [Davis, 1985]. They are influenced by motions at their own scales, of the order of 1 m in the horizontal, and communicate their position at intervals of the order

23 of minutes for typical coastal applications. Measurements from HF radar and
24 CODE type drifters are considered compatible in the vertical since they both
25 sample approximately the first meter of water below the surface [Stewart and
26 Joy, 1974]. Regarding instrument errors, drifters are subject to windage and
27 slippage, that for CODE drifters are estimated to be within 1-3 cm/s for winds
28 up to 10 m/s [Poulain et al., 2009]. For HF radar measurements, uncertain-
29 ties are more complex and less easily quantifiable [Chapman and Graber, 1997].
30 They can be due to actual measurement errors related to uncalibrated antenna
31 patterns or radio interferences [Gurgel and Barbin, 2008], but they can also be
32 due to sea state or dependence from the methods used to reconstruct the vector
33 velocity from the radial velocities measured by the antennas [Kohut and Glenn,
34 2003; Kohut et al., 2012]. Comparisons between velocity measurements from
35 HF radars and drifters indicate typical differences of the order of 5-15 cm/s
36 [Chapman et al., 1997; Emery et al., 2004; Essen et al., 2000; Kaplan et al.,
37 2005; Paduan and Rosenfeld, 1996; Rypina et al., 2014; Shay et al., 1998a,b,
38 2001, 2007]. These differences can be due not only to the measurement errors of
39 the two platforms, but also to the different nature of the measurements, i.e. the
40 fact that HF radar velocities are averaged over cells of more than 1 km and over
41 time intervals of the order of 1 h [Paduan et al., 2006; Paduan and Washburn,
42 2013], while drifters provide more localized spatial and temporal information.
43 Indeed, in many cases it has been shown that the differences are within the
44 expected environmental variability inside the radar averaging cells, such as in
45 Mantovanelli et al. [2011], Ohlmann et al. [2007].

46 Another very important tool used to estimate Lagrangian transport is the use
47 of numerical models. The accuracy of numerical models has greatly improved
48 in the last years, in terms of resolution, data assimilation and use of specific dy-
49 namical system methods to compute transport [Haller and Poje, 1998; Olascoaga
50 et al., 2006]. In particular, open ocean models assimilating altimetric data ap-
51 pear able to capture several transport features induced by large and mesoscale
52 structures [Olascoaga and Haller, 2012]. Capturing smaller scale coastal cur-
53 rents and their effects on transport can be more challenging for models, unless

54 in regions with long standing observatories [Haza et al., 2007; Kuang et al.,
55 2012] where models have been thoroughly validated and often local in situ data
56 are assimilated.

57 Lagrangian data from surface drifters or subsurface profiling floats are ex-
58 pected to be particularly useful for assimilation in models to improve estimates
59 of Lagrangian transport. Various methods have been developed for Lagrangian
60 data assimilation in the last decades [Krause and Restrepo, 2009; Kuznetsov
61 et al., 2009; Molcard et al., 2003], and they have been tested with positive re-
62 sults especially for subsurface floats in open ocean [Taillandier et al., 2006b].
63 Assimilation of surface drifters especially in coastal regions is still challenging,
64 since drifters provide information on the upper ocean which is highly influenced
65 by air-sea interaction and submesoscale processes that are often only partially
66 resolved by models. Also, the high deviation from geostrophy that is likely to
67 dominate surface dynamics poses additional problems on the balancing of the
68 other model state variables for assimilation. For these reasons, surface drifters
69 have been used in a number of applications to blend rather than truly assimilate
70 information in models [Chang et al., 2011; Taillandier et al., 2006a]. Blending
71 corrects the surface velocity field when drifter data are available but it does not
72 provide a complete corrected model state, so that re-initialization cannot be
73 performed and full forecast capabilities are not provided. On the other hand,
74 blending has the advantage that it can be quickly applied to any available veloc-
75 ity field and it can be very useful in many practical and operational applications
76 for nowcast.

77 In particular, the LAgrangian Variational Analysis (LAVA) is a method that
78 has been used for both, full assimilation of subsurface floats at 350 m [Taillandier
79 et al., 2006b, 2010] and blending of surface drifters in various and different re-
80 gions such as the Adriatic Sea (Mediterranean Sea) and the Kuroshio extension
81 [Chang et al., 2011; Taillandier et al., 2008]. LAVA combines information from
82 Lagrangian instruments with model velocity fields requiring that the distance
83 between observed trajectories and synthetic trajectories computed from the ve-
84 locity field is minimized. The method is therefore strictly Lagrangian, in the

85 sense that it directly uses the Lagrangian position information rather than the
86 “pseudo-Eulerian” velocity computed from the trajectories, and its correction
87 is specifically aimed at optimizing Lagrangian transport. In principle, LAVA
88 can be used not only to correct velocity fields from models but also from any
89 other data source, like for example from HF radars. When other information
90 are not available, the reconstruction of the velocity field can also be obtained
91 with LAVA using Lagrangian data only [Taillandier et al., 2006a].

92 In this paper we use LAVA to blend drifter data with velocity fields from HF
93 radars and models, as well as to reconstruct velocity from drifters only, in the
94 French coastal area of the Ligurian Sea in front of Toulon, in the North West-
95 ern Mediterranean Sea (Fig.1). While blending with model results has been
96 tested in a number of previous applications [Chang et al., 2011; Taillandier
97 et al., 2008], blending with HF radar is performed here for the first time, and
98 it deserves a brief discussion to clarify its interpretation. The process of blend-
99 ing two data sets, each one with its peculiarity and uncertainties, is intended
100 to provide an optimized field that takes into account the best aspects of each
101 platform. Drifters provide local information with relatively high precision while
102 HF radars provide extensive spatial information. The idea is to use the local
103 information from the drifters not only to correct possible errors in the radar field
104 (occurring for instance in case of difficult sea state, see Essen et al. [2000], or
105 unknown antenna pattern, see Kohut and Glenn [2003]), but also and foremost
106 to re-establish details of the environmental space variability that are smoothed
107 by radar averaging. In case of strong horizontal shear, for instance, it can be ex-
108 pected that drifter information can help sharpening the velocity field increasing
109 the gradients. This potential use of blending different data sets is investigated
110 here, focusing on Lagrangian transport.

111 The goal of the work is twofold: (i) to investigate the characteristics and
112 predictability of Lagrangian transport in a coastal area; (ii) to provide in-
113 formation for practical applications involving rapid response in case of ac-
114 cidents at sea. The work is performed in the framework of the EU-MED
115 project TOSCA (Tracking Oil Spills and Coastal Awareness network, [http:](http://)

116 //www.tosca-med.eu), that is aimed at investigating and testing science-based
117 methodologies, best practices, and response plans in case of accidents at sea.
118 In particular the TOSCA project consists in the development of a coastal
119 monitoring and forecasting network based on HF radars and new generation
120 drifting instruments and models, aimed at optimizing the response of local au-
121 thorities to marine accidents, with a special emphasis on oil spill pollution and
122 on SAR operations. Under this frame, the results of the present work can be
123 useful to indicate an optimized use of HF radar and drifter data to improve
124 estimates of Lagrangian transport.

125 The circulation in the study area (Fig.1(a)) was investigated with a dedi-
126 cated experiment during the month of August 2012, targeting a region of the
127 order of 50 km range, covered by a high resolution coastal model (GLAZUR64,
128 Ourmières et al. [2011]) and by an active HF radar installation. During the ex-
129 periment, a total number of 20 drifters (including redeployments) were launched,
130 and a number of LAVA blending configurations are tested here. A first set of
131 LAVA experiments is carried out blending the trajectories of $N_{Dft} = 7$ drifters,
132 while the remaining 13 trajectories are used as testing for the results (“control”
133 trajectories). To assess sensitivity to the number of drifters used in the LAVA
134 blending, a second set of experiments is then performed downgrading N_{Dft} to
135 5. The LAVA algorithm is applied to both HF radar and model velocity fields,
136 and the results of the blended fields are compared to the results of the original
137 fields in terms of Lagrangian transport. In addition to the hindcast/nowcast
138 capability of the system, we also test a very simplified version of “forecast”,
139 where the velocity is assumed to have a temporal persistency of the order of
140 hours to a day. In other words, the velocity field at a given time is maintained
141 steady over a certain period and it is used to compute Lagrangian transport.
142 The method implies of course a simplified representation of the current field, but
143 its testing is relevant for operational situations, to verify whether or not some
144 guidance can be provided at least for the first few hours. It is important to keep
145 in mind that the Mediterranean Sea has small tidal ranges being connected to
146 the Atlantic Ocean through a narrow entrance (the Gibraltar Strait) [Arabelos

147 et al., 2011; Pugh, 1987]. For this reason, relatively slowly varying mesoscale
148 flows are expected to be more relevant than tidal fluctuations.

149 The Toulon experiment and the data from HF radars, drifters and models
150 are described in Section 2. The description of the LAVA method is provided in
151 Section 3 together with the description of the main diagnostics used to quantify
152 the results. The LAVA application to HF radar and model velocities is shown
153 in Section 4 in terms of Eulerian statistics and Lagrangian transport estimates.
154 A summary and concluding remarks are given in Section 5.

155 **2. The TOSCA-Toulon experiment and data sets**

156 The TOSCA-Toulon experiment took place during the month of August
157 2012 in the area shown in Fig.1. The circulation in the area is part of the
158 Ligurian/Northern Current system, that flows cyclonically along the coasts of
159 Italy and France. The current is characterized by high variability at many scales,
160 from seasonal to mesoscale and submesoscale. It is especially energetic during
161 winter, with a well defined core at 10-40 km from the coast, while during summer
162 it is weaker and often spreading offshore [Alb erola et al., 1995a]. Mesoscale
163 structures are also less energetic in summer, even though present all year round
164 and characterized by meandering activities with periods in two main bands of
165 approximately 3-6 and 10 days respectively [Sammari et al., 1995]. The typical
166 Rossby radius R_d is of the order of 10 km, even though smaller structures of
167 the order of 5 km can be found [Marullo et al., 1985]. In particular the region
168 of Toulon is characterized by very high variability, with frequent open ocean
169 intrusions and formation of jets and eddies [Bellomo et al., 2013; Bosse et al.,
170 2013; Guihou et al., 2013].

171 During the period August 5-10, drifters were repeatedly launched from the
172 R/V Urania. A HF radar system was operative and a high resolution model
173 was running in real time. Drifter trajectories, model and radar coverages are
174 all shown in Fig.1. Details on their data are given in the following.

175 *2.1. HF radar data*

176 The HF radar installation is based on the WERA technology [Gurgel et al.,
177 1999a] and relies on two systems. The first one (Fort Peyras, “FP” in Fig.1)
178 has a quasi-monostatic configuration with an irregular, W-shaped 8-antenna
179 receiving array and 2 monopoles performing the emission while forming a zero
180 in the direction of the receiver. The peculiarity of the receiving array geometry
181 is imposed by the environment of the site, a dismissed military base. The
182 second system has a bistatic configuration, with the single emitter antenna
183 (Porquerolles island, “P” in Fig.1) located at about 17 km from the receiver
184 (Cap Bénat, “CB” in Fig.1), a regular linear 8-antenna array.

185 The two systems operate at a frequency of 16.1 MHz with bandwidth of
186 50 kHz, giving a range resolution in the radial direction of 3 km. Antenna
187 patterns are routinely measured almost every year. The azimuthal processing
188 is done with the MUSIC (MUltiple SIgnal Characterization) direction finding
189 algorithm employed routinely in CODAR [Lipa et al., 2006] and less frequently
190 in WERA systems [Molcard et al., 2009; Sentchev et al., 2013], with a nominal
191 2° resolution. Current maps are produced every 20 min by integrating over the
192 previous hour. As evidenced by Kohut et al. [2012], the vector computation
193 accuracy partly depends on the algorithm used to compute the velocity field
194 as well as on the Geometric Dilution Of Precision (GDOP). In this case, total
195 velocities are computed on a regular 2 km grid with a local interpolation method
196 which, at each cartesian grid point, minimizes the Mean Square Error (MSE)
197 between the projection of the cartesian velocity onto the radial directions and
198 the radial velocities available within a circle with radius 3 km [Lipa and Barrick,
199 1983]. To reduce errors from GDOP [Chapman et al., 1997]), totals are only
200 computed when the angle between radial data from the two sites is within the
201 range $30\text{-}150^\circ$, which corresponds to GDOP values smaller than 2.5.

202 It has been shown [Stewart and Joy, 1974] that HF radars retrieve cur-
203 rent velocities which are vertically averaged through an exponential weighting
204 function with a characteristic depth $\lambda_w/4\pi$, where λ_w is the wavelength of the
205 Bragg-resonant sea waves. Since for a monostatic HF radar and approximately

206 for a bistatic one, too, $\lambda_w = 0.5\lambda_0$, with λ_0 the electromagnetic wavelength, at
207 the frequency at which the system was operated in our case $\lambda_w/4\pi$ gives ~ 75
208 cm for the equivalent depth.

209 *2.2. Model data*

210 The model configuration (GLAZUR64) is based on the primitive equation
211 ocean circulation model NEMO [Madec, 2008]. The resolution of GLAZUR64
212 is uniform in the horizontal direction and set to $1/64^\circ$ (about 1.5 km). In
213 the vertical direction, 130 z-levels are considered with the first level at 0.5 m
214 depth and decreasing resolution from 1 m near the surface to 30 m near the sea
215 bed. This high-resolution three-dimensional mesh allows to well reproduce the
216 mesoscale processes.

217 The model domain is shown in Fig.1 and covers the Gulf of Lions and part
218 of the Western Ligurian Sea, featuring two open boundaries, at East and South.
219 Initial and boundary conditions for temperature, salinity and velocity fields
220 are provided by a large scale operational model (MERCATOR OCEAN, <http://www.mercator-ocean.fr>).
221 MERCATOR products include assimilation of
222 satellite data (namely, sea surface temperatures and surface level anomalies)
223 and in situ hydrographic profiles, providing daily averaged oceanic fields on
224 a $1/12^\circ$ horizontal resolution with 50 vertical levels. Surface conditions rely
225 on the atmospheric data from the Meteo-France operational regional model
226 ARPEGE (spatial and temporal resolutions of 10 km and 3 h, respectively).
227 Such a high spatio-temporal resolution is essential to correctly reproduce wind
228 induced oceanic features and variability in the area, especially when the focus
229 is on the surface circulation [Madec, 2008; Schaeffer et al., 2011].

230 Bulk formulation is used and requires zonal and meridional wind components
231 at 10 m, temperature and specific humidity at 2 m, precipitation, radiative
232 and solar fluxes. Atmospheric forcings are linearly interpolated in time. A
233 bicubic spatial interpolation is achieved for wind forcing, while spatial linear
234 interpolation is used for the other atmospheric forcing.

235 The GLAZUR64 model has been evaluated [Ourmières et al., 2011; Guihou

236 et al., 2013] by comparing the simulations with in situ measurements and remote
237 sensing data and it was shown to realistically simulate the Northern Current
238 mesoscale variability at sub-regional scale.

239 The simulation considered here starts in June 2011 from initial conditions
240 provided by the large scale operational model, allowing a sufficient spin-up ad-
241 justment prior to the period of interest (see Guihou et al. [2013] for a validation
242 study of GLAZUR64).

243 Two-dimensional surface velocity fields used in the present analysis have a
244 time resolution of 1 h.

245 *2.3. Drifter data*

246 During the Toulon experiment, 20 CODE drifters (including redeployments)
247 were launched. CODE drifters [Davis, 1985; Poulain, 1999] measure the current
248 in the first meter under the sea surface, and they have been chosen for the
249 TOSCA experiment since they provide information on surface coastal circulation
250 that is relevant for the targeted applications. Also, CODE data are expected
251 to be compatible with HF radar data, because the equivalent integration depth
252 is ~ 75 cm. The drifters were equipped with Global Positioning System (GPS)
253 receivers with an accuracy of approximately 5-10 m. Drifter positions, retrieved
254 every 15 min, were edited to remove spikes and offsets and interpolated at
255 uniform intervals. As observed by Edwards et al. [2006] currents measured by
256 drifters can be partially affected by other forcings, such as wind stress and
257 slippage. In the specific case of CODE drifters, the comparison with current
258 meter measurements [Davis, 1985; Poulain et al., 2009] showed that drifters
259 follow surface currents to within 3 cm/s with wind conditions up to 10 m/s.
260 During the analyzed time windows, the wind never reached a speed comparable
261 to this value (not shown). Moreover, Poulain et al. [2002] performed specific
262 slippage measurements with acoustic current meters positioned at the top and
263 at the bottom of the CODE (about 1 m apart along the vertical) and showed
264 that drifters follow surface currents within 2 cm/s and that they are consistent
265 with the near-surface Ekman dynamics.

266 During the Toulon experiment the drifters were launched in triplets and pairs
267 characterized by an initial distance of the order of 100 m, while the typical rel-
268 ative distance between triplets was of the order of a few km. This deployment
269 strategy is chosen to evidence the separation among clustered drifters due to
270 the small scale variability and its effect on the Lagrangian transport. In fact,
271 the experiment has a multi-scale aim: it assures a quasi-homogeneous coverage
272 at mesoscale while providing also small scale information, as specifically inves-
273 tigated in Section 4.3. The overall drifter distribution is shown in Fig.1, with
274 color-coded trajectories indicating time evolution. A first set of 14 trajectories
275 was launched on August 5 (blue color). This set drifted mostly westward fol-
276 lowing the current and exited the radar region by August 7. Later, 6 of these
277 drifters were picked up and redeployed in the HF radar zone on August 8 (yellow
278 color).

279 **3. Methodology: LAVA**

280 LAVA is a variational method developed by Taillandier et al. [2006a] and ap-
281 plied in a number of cases for both full assimilation of subsurface float data and
282 blending of surface drifters [Chang et al., 2011; Taillandier et al., 2006b, 2008].
283 The basic concept behind LAVA is that Lagrangian data are used to correct
284 an available Eulerian velocity field, used as first guess, by minimizing the dis-
285 tance between the observed positions and the positions of numerical trajectories
286 advected in the Eulerian field. The correction along the trajectories is spread
287 using a diffusion equation [Derber and Rosati, 1989; Weaver and Courtier, 2001]
288 with space scale R . The procedure is applied sequentially with individual time
289 sequences of length T_a , given by $T_a = m\Delta t$, where m is an integer and Δt is
290 the time step over which data are provided.

291 The underlying hypothesis is that T_a is significantly shorter than the persis-
292 tency time of the Eulerian velocity, T_E , and also of the Lagrangian time scale
293 of the drifter, T_L : $T_a < T_E$, $T_a < T_L$. The space scale R , on the other hand,
294 is assumed to be of the order of the typical Rossby Radius R_d in the area,

295 and is of course greater than the grid size Δx over which the Eulerian field is
296 implemented.

297 We also notice that drifter trajectories used in LAVA are required to have
298 a relative distance greater than $\approx 2\Delta x$, in order to avoid conflicting velocity
299 information at the grid scale in the blending process. For this reason, given that
300 the TOSCA experiment drifters have been launched in close triplets or pairs, the
301 original 20 trajectories need to be subsampled for LAVA application. The maxi-
302 mum number of drifters obeying the requirement is $N_{Dft} = 7$, corresponding to
303 choosing a single trajectory in each drifter group (see Fig.2). In the following,
304 the remaining 13 trajectories are considered as control data and are used to test
305 the LAVA results. In addition to the main configuration with $N_{Dft} = 7$, tests
306 have also been performed downgrading the number of drifters blended in LAVA
307 to $N_{Dft} = 5$, simply discarding some of the trajectories. In order to charac-
308 terize each drifter configuration, a parameter $\overline{d_L}$ is introduced, that measures
309 the mean initial distance between drifters used in the LAVA blending and the
310 control ones used to test the results. The meaning and relevance of this param-
311 eter are discussed in Section 3.1.2, with reference to the introduced Lagrangian
312 metrics.

313 *3.1. Diagnostics of the results: Eulerian and Lagrangian metrics*

314 The results presented in Section 4 are diagnosed using the Eulerian and
315 Lagrangian metrics detailed below.

316 *3.1.1. Eulerian metrics*

317 The Eulerian metrics are designed to provide a quantitative estimate of the
318 difference (correction C) between the original and LAVA blended velocity fields.
319 These corrections are interpreted as an indication of the good agreement between
320 the Eulerian velocity field and the in situ drifter data.

321 In Section 4, we first present a visual comparison of daily averages of the
322 original and LAVA blended fields and of the amplitude of the vector difference
323 between the two

$$C_{daily} = \sqrt{(\langle u_{or} \rangle - \langle u_{LA} \rangle)^2 + (\langle v_{or} \rangle - \langle v_{LA} \rangle)^2}, \quad (1)$$

324 where u and v are the two velocity components, $\langle \rangle$ indicates daily averages
 325 and the subscripts u_{or}, u_{LA} indicate the original and LAVA blended velocity
 326 fields respectively.

327 We then provide a bulk assessment of the correction versus time, computing
 328 for each time step a spatial average value of the correction normalized by the
 329 original average value:

$$C_{norm}(t) = C_{int}(t)/U_{or}(t), \quad (2)$$

where

$$C_{int} = \langle \sqrt{(u_{or} - u_{LA})^2 + (v_{or} - v_{LA})^2} \rangle_A, \quad (3)$$

$$U_{or} = \langle \sqrt{u_{or}^2 + v_{or}^2} \rangle_A .$$

330 To preserve the contribution of all drifters, at each time t , the spatial average
 331 $\langle \rangle_A$ is made over the regions where the kinetic energy of the corrected fields
 332 is at least 30% of the least energetic drifter correction. Sensitivity tests have
 333 been performed in the threshold range 10 – 50% and the results are robust (not
 334 shown). The same kinetic energy criterion, applied to define the region affected
 335 by corrections, is used to perform weighted temporal averages of the velocity
 336 field reconstructed from drifters only in order to avoid biases due to the lack of
 337 velocity field estimates where drifters are not present.

338 3.1.2. Lagrangian metrics

339 Lagrangian diagnostics are designed to measure and compare the capability
 340 to estimate Lagrangian transport of original and blended velocity fields. Be-
 341 cause Lagrangian predictability in the upper ocean is expected to be in the
 342 order of 1 day and the same time window is relevant for practical operational
 343 applications, observed trajectories are first split in segments of 24 h. Only
 344 trajectories at least 24-h long are considered and each segment is treated as a
 345 different observed trajectory in the statistics.

346 We first compute a simple estimate of the transport uncertainty, $D(t)$, given
347 by the average distance at each time between the control (i.e. not used in
348 LAVA) observed trajectories and those obtained numerically using the original
349 and LAVA blended velocities. Specifically, numerical drifters are initialized at
350 the same positions as the observed control drifters and are obtained integrating
351 the Eulerian velocity in time with a fourth-order Runge Kutta scheme. Thus,
352 $D(t)$ is defined as

$$D(t) = \langle \sqrt{(x_d - x_n)^2 + (y_d - y_n)^2} \rangle_d, \quad (4)$$

353 where (x, y) are the components of the drifter position at time t and the sub-
354 scripts d and n indicate in situ and numerical drifters, respectively. The average
355 $\langle \rangle_d$ is performed over the number of control drifters and calculated over a time
356 period of 24 h.

357 At each time, $D(t)$ is compared to $D_0(t)$, i.e. the average absolute dispersion
358 of the observed drifters over a time period of 24 h, defined as

$$D_0(t) = \langle \sqrt{(x_d - X_d)^2 + (y_d - Y_d)^2} \rangle_d, \quad (5)$$

359 where (X_d, Y_d) are the components of the initial positions of the drifters. D_0
360 provides a measure of the average distance covered by drifters at each time t and
361 is commonly referred to as “persistency error”, i.e. the error that corresponds
362 to a zero prior knowledge, i.e. assuming that particles do not move from their
363 initial conditions [Ullman et al., 2006]. Practically, any $D(t)$ smaller than $D_0(t)$
364 represents an improvement with respect to a zero prior knowledge.

365 The metric $D(t)$ is computed averaging over control drifters that are not used
366 in the LAVA blending and is expected to be influenced by the initial distance
367 d_L between the control and blended drifters. We remark that since LAVA
368 correction is spread over a range within R from the blended drifters, the LAVA
369 assessment must be performed using control drifters within the same range R .
370 Conceptually, we can expect that the LAVA correction is most effective near
371 the blended trajectories, while it decreases at increasing distance from them
372 and it vanishes at distances greater than R . As a consequence, the uncertainty

373 estimate $D(t)$ is expected to increase at increasing values of the mean d_L . The
 374 details of this dependence might not be trivial and might be relevant for practical
 375 applications. A very small $\overline{d_L}$, of the order of 0.1 – 1 km, corresponds to an
 376 operational situation where drifters are launched relatively close to the quantity
 377 to be monitored, for instance in case of an oil spill or SAR episode timely
 378 observed and reported. A larger $\overline{d_L}$ of a few km, on the other hand, could
 379 correspond to a case when the initial conditions of the accident are poorly
 380 known, and a large scale coverage is performed. In Section 4, we provide some
 381 insights on the dependence from $\overline{d_L}$ considering LAVA experiments with $N_{Dft} =$
 382 7 and $N_{Dft} = 5$ drifters, that correspond to different $\overline{d_L}$. Given the overall small
 383 number of drifters (20 total) and their coverage, the number of configurations
 384 is necessarily reduced and testing is necessarily limited. Nevertheless, some
 385 indications can be drawn from the results.

386 3.2. LAVA experiments

387 In this study, 7 different LAVA experiments are configured (see Tab.1), vary-
 388 ing the first guess Eulerian velocity fields and the parameters N_{Dft} , Δx , Δt ,
 389 T_a and $\overline{d_L}$. In all the cases, R is set to 7 km, i.e. of the order of the estimated
 390 Rossby Radius R_d in the area [Marullo et al., 1985; Robinson et al., 2001].

391 In the first 3 experiments, E_R , E_M , E_D , the maximum number of blended
 392 trajectories is used, i.e. $N_{Dft} = 7$. The 7 drifters are shown in Fig.2, and
 393 they correspond to 5 drifters providing a good coverage of the area during a
 394 first period of ≈ 2 days, 5 – 6 August, and other 2 situated mostly in the
 395 southern part of the domain in a second period during the last 2 days, 8 – 10
 396 August. A data gap occurs in the intermediate period around August 7. The
 397 corresponding parameter $\overline{d_L}$ for this drifter configuration is approximately 250
 398 m (Tab.1), with slight value differences among the experiments because of the
 399 different discretizations of the first guess velocity, as explained below.

400 The parameters, Δx , Δt , T_a are chosen to retain in LAVA the highest spatial
 401 and time resolutions possible, still allowing for the optimal minimization of the
 402 misfits between positions. In the first experiment, E_R , the first guess Eulerian

403 velocity is given by the radar, and grid size and time step are maintained as in
 404 the original fields: $\Delta x = 2$ km and $\Delta t = 20$ min. The same strategy is followed
 405 for E_M , where the first guess is given by the model fields with Δx and Δt set
 406 to $1/64^\circ$ and 1 h respectively. Since radar data are more frequent in time, the
 407 analysis time scale T_a for the minimization in E_R can be chosen shorter (2 h)
 408 than for E_M (4 h).

409 In the third experiment, E_D , LAVA is applied with zero first guess, i.e.
 410 assuming no prior information available from models or radars, so that LAVA
 411 provides a reconstruction of the velocity based only on the positions of the
 412 drifters. For this case, LAVA parameters are set as for E_R , namely $\Delta x = 2$ km,
 413 $\Delta t = 20$ min and $T_a = 2$ h.

414 In the other four experiments, (E_R^{T1} E_R^{T2} E_M^{T1} E_M^{T2} , Tab.1), the sensitivity to
 415 the number of drifters used in the LAVA blending N_{Dft} , and to the configuration
 416 parameter $\overline{d_L}$ is assessed. Thus, N_{Dft} is downgraded to 5 and two $\overline{d_L}$ values are
 417 considered, namely $\overline{d_L} \approx R/4$ and $\overline{d_L} \approx R/2$. The disregarded drifters belong to
 418 the pairs/triplets around $42^\circ 39'N$ $6^\circ 18'E$, $42^\circ 46'N$ $6^\circ 19'E$ and $42^\circ 48'N$ $6^\circ 17'E$,
 419 respectively (Figs. 2(a) and 2(b)). In the two experiments E_R^{T1} and E_R^{T2} (E_M^{T1}
 420 and E_M^{T2}) this sensitivity is explored using the first guess Eulerian velocities
 421 from the radar (model).

422 In all the experiments, T_a allows to resolve the local inertial period ($T_I \sim 17$
 423 h, Picco et al. [2010]), while original drifter positions are interpolated according
 424 to the different temporal resolutions of the velocity field in the experiments.
 425 Tab.1 summarizes all parameters for the 7 experiments.

426 4. Results

427 Results illustrate the impact of applying LAVA to the original velocity fields
 428 from radar and model. Results from the first 3 experiments, E_R , E_M , E_D ,
 429 with $N_{Dft} = 7$ drifter trajectories are discussed first, in terms of Eulerian and
 430 Lagrangian metrics (Sections 4.1 and 4.2). The sensitivity of the results is then
 431 investigated via the experiments with only $N_{Dft} = 5$ drifters (Section 4.3).

432 Finally a discussion on the forecasting skills of the blended fields is carried out
433 in Section 4.4.

434 4.1. Eulerian statistics

435 Examples of daily average fields are shown in Figs.2, 3 and 4 where we
436 compare original data with those obtained from LAVA experiments E_R , E_M
437 and E_D respectively. The evolution of the normalized correction C_{norm} (eq. 2)
438 is shown in Fig.5.

439 Two examples of original radar and model daily averaged fields for days
440 August 5 and 9 are shown in the upper panels of Fig.2 and 3 respectively.
441 Arrows indicate velocity vectors, colors indicate velocity amplitudes, and black
442 lines are for the drifter trajectories used in LAVA during the considered periods.
443 In the middle panels of the same figures instead, arrows indicate LAVA blended
444 fields while colors indicate C_{daily} (eq. 1) correction amplitudes. Finally in the
445 lower panels, the correction vectors are shown, i.e. the differences between the
446 LAVA blended and original velocities.

447 We start by discussing the radar daily averages in Fig.2. The original radar
448 velocity (upper panels) depicts a well defined and persistent boundary current.
449 The current is approximately zonal and flows westward along the coast (north of
450 $\sim 42^\circ 50'N$) with velocities up to 50 cm/s. The southern region is more variable
451 and less energetic, with velocities less than 20 cm/s. Inspection of similar figures
452 at different times t (not shown) reveals the high time variability of this area,
453 suggestive of meanders and recirculation. A strong inertial signal (with period
454 of ~ 17 h) is also evident, as suggested also by drifter trajectories.

455 The LAVA blended fields from experiment E_R (middle panels) show very
456 similar patterns with respect to the originals, indicating that drifter motion is
457 in good agreement with the radar velocity. This is quantitatively shown by
458 the C_{daily} values. During the first period (left panel), C_{daily} reaches maximum
459 values of the order of 10 cm/s in the upper northern region where the veloc-
460 ities are highest, while during the second period (right panel) the corrections
461 are significantly lower than 10 cm/s and limited to the less energetic southern

462 region. The velocity corrections (lower panels) show a clear pattern of shear
463 enhancement in the coastal jet during the first period.

464 The original model results without LAVA (Fig.3, upper panels) do not show
465 the presence of the zonal boundary current in the northern region. Rather, the
466 whole area is characterized by a strong meandering activity, with meridional
467 velocities both in the inshore and offshore directions, with highest velocities of
468 the order of 50 cm/s.

469 The LAVA blended fields from experiment E_M (middle panels) and the cor-
470 rection vectors (lower panels) present significantly different patterns from the
471 original fields especially during the first period (left panel). In particular, the
472 zonal boundary current can be now seen in the northern part of the domain
473 during the first period, due to the presence of the drifters, and the correction
474 reaches significantly higher values than for the radar, up to 30 cm/s. This is
475 less evident in the second period (right panel), when the drifters cover only the
476 southern region and the correction is more limited and similar to the one of
477 the radar in magnitude. We notice that the velocity pattern depicted by the
478 model during the first period with a high meridional shear can be indicative of
479 meandering and jet propagations, as previously observed in this area [Guihou
480 et al., 2013]. It is possible that the model results depict such pattern instead of
481 a zonal current because of a time error or shift in the propagation of a pertur-
482 bation, as it can happen quite commonly especially for nonlinear coastal flows
483 in non-assimilating models.

484 It is interesting to compare radar and model velocities with results obtained
485 in the experiment E_D , i.e. reconstructing the velocity fields from drifters only.
486 The daily fields in Fig.4 show patterns highly consistent with the radar results
487 in Fig.2, especially during the first period (left panel) when the drifter coverage
488 extends throughout the domain. In the second period (right panel) the coverage
489 is limited to the southern part and shows a good qualitative agreement with both
490 radar and model fields. These results indicate that even when radar and model
491 information are not available, the use of drifters only can lead to significant
492 information in terms of velocity fields, of course if an adequate enough coverage

493 is provided.

494 Time series of C_{norm} (eq.2) for E_R and E_M are shown in Fig.5. During
495 the first period, with extended drifter coverage, C_{norm} is significantly higher
496 for the model than for the radar reaching values of almost 70%, consistently
497 with what shown by the daily results in Figs.2 and 3. The average correction
498 (Tab.2) is almost double for the model with respect to the radar (36% and 19%
499 respectively). During the second period, instead, when the coverage is limited
500 to the southern region, model and radar corrections are very similar and the
501 average is actually slightly higher for the radar (25% and 22% for E_R and E_M ,
502 respectively).

503 4.2. Lagrangian statistics

504 In order to visualize the effects of LAVA on Lagrangian transport, we first
505 show in Fig.6 some qualitative examples of trajectories computed from original
506 and LAVA blended radar and model velocities, comparing them with in situ
507 drifter trajectories. The radar results (upper panels) show that the numerical
508 trajectories computed from the original fields (green lines in left panel) are
509 relatively similar to the drifter ones (black lines). The radar performance is
510 further improved by the LAVA blending in the experiment E_R (purple lines in
511 right panel). Model results (lower panels) are more striking, with trajectories
512 from the original fields (left panel) considerably different from the drifters and
513 in some cases diverging in opposite directions. When LAVA is applied in the
514 experiment E_M (right panel) though, results are improved, with trajectories
515 very close to the drifters and comparable to the radar results. This suggests
516 that LAVA represents an effective method to enhance trajectory estimates and
517 therefore also Lagrangian transport analysis.

518 A quantitative measure of the LAVA effects on transport estimates is given
519 by the statistical quantities $D(t)$ and $D_0(t)$, (eqs.4 and 5), respectively. Results
520 are shown in Fig.7 (left panels) for radar (upper panel) and model (lower panel).
521 The dotted line indicates the absolute drifter dispersion $D_0(t)$ considered as a
522 reference, reaching values of $\approx 12-13$ km after 24 h. The radar results indicate

523 that the trajectory uncertainty $D(t)$ computed using the original fields (pink
524 line) grows almost linearly, reaching $\approx 6-7$ km at 24 h, i.e. approximately half
525 of $D_0(t)$. For LAVA blended E_R fields (green line), the uncertainty decreases
526 with a maximum of 2 km at 24 h. For the model, $D(t)$ of the original field (red
527 line) is of the same order as $D_0(t)$, actually reaching slightly higher values of
528 14 km at 24 h. The LAVA correction for the E_M experiment induces a striking
529 uncertainty decrease, with maximum values less than 2 km, i.e. of the same
530 order or smaller than the radar ones.

531 These results show the great advantage of using LAVA to improve trans-
532 port estimates, and they suggest possible practical consequences for operational
533 problems such as SAR or identification of pollutant spreading. In cases when
534 no information from radar or model or other sources are available, it can be
535 expected that a search range will be of the order of the uncertainty on particle
536 positions quantified by the absolute dispersion $D_0(t)$, i.e. of the order of 10-15
537 km after 24 h in this area. When radar information are available, the results
538 in Fig.7 suggest that the uncertainty associated with radar based trajectories
539 $D(t)$ is approximately half than $D_0(t)$, so that the range can be decreased to
540 approximately 6 km after 24 h. For model information, instead, at least in the
541 case we considered, the uncertainty is of the same order of $D_0(t)$ so that the
542 range cannot be decreased. When LAVA is applied, though, the uncertainty
543 after 24 h is decreased to only 2 km for both radar and model, suggesting a
544 drastic reduction in the search range.

545 We notice that in the case of the model, the errors corrected by LAVA are
546 likely to be due to the nonlinear propagation of meanders that are not correctly
547 depicted by the simulation. For the HF radar, instead, regardless the accuracy
548 of the radar settings and the data processing algorithms, the LAVA blending is
549 expected to improve the results because the localized drifter data restore part of
550 the environmental variability smoothed by the radar. This is shown for instance
551 by the enhanced shear correction in Fig.2 (lower panels).

552 We recall that all these results are related to hindcast or nowcast applica-
553 tions, while a discussion on possible forecast applications will be provided in

554 Section 4.4.

555 *4.3. Sensitivity tests using downgraded configurations*

556 All the results discussed in Sections 4.1 and 4.2 concern the first 3 exper-
557 iments which blend 7 trajectories and use the remaining 13 to evaluate $D(t)$.
558 In these experiments, $\overline{d_L}$ is approximately 250 m. Here we test the sensitivity
559 of the $D(t)$ results considering 2 different realizations of LAVA with a reduced
560 number of drifters, $N_{Dft} = 5$, and different values of $\overline{d_L}$, $\approx R/4$ and $\approx R/2$
561 respectively, while maintaining the same drifters to evaluate $D(t)$ (see Tab.1).

562 Results in terms of $D(t)$ are shown in Fig.7 (right panels), for radar (E_R^{T1}
563 and E_R^{T2} , upper panel) and model (E_M^{T1} and E_M^{T2} , lower panel). The depen-
564 dence on $\overline{d_L}$ can be easily seen, comparing also with the results of the main
565 LAVA application (left panels). $D(t)$ generally increases with $\overline{d_L}$, even though
566 at different rates between radar and model. For the radar (Fig.7b), the results
567 of E_R^{T1} are similar to E_R with D reaching ≈ 2 km after 24 h, while there is a
568 clear increase to ≈ 4 km for E_R^{T2} . For the model (Fig.7d), an increase to ≈ 3 km
569 is already visible for E_M^{T1} while in E_M^{T2} values greater than 4 km are reached.
570 These differences between radar and model are probably only marginally sig-
571 nificant, given the size of standard deviations especially for the model, but it is
572 nevertheless to be expected that the model results are more sensitive given that
573 the LAVA correction is significantly greater than for the radar.

574 In summary, the results confirm that $D(t)$ increase with $\overline{d_L}$, but also suggest
575 that the growth is relatively limited at least in the considered range, with $\overline{d_L}$
576 smaller than the Rossby Radius R_d . We recall in fact that R_d , in the area, is of
577 the order of 5-10 km, comparable to the LAVA space scale R , which is set to 7
578 km in all experiments. Maximum values of d_L are of the order of R (not shown)
579 and, as a consequence, the LAVA correction is still expected to be significant,
580 and the $D(t)$ results only partially downgraded. For instance, if we compare
581 the highest values of $D(t)$ obtained for the E_M^{T2} experiment with $D_0(t)$ and the
582 original $D(t)$ (lower left panel), we see that the E_M^{T2} values are approximately a
583 third of the original ones indicating a persisting significant advantage in using

584 LAVA. Of course if d_L values were greater than R_d , we could expect that the
585 advantage of using LAVA will cease.

586 From the application point of view, the results indicate that even when the
587 initial conditions of an accident are only approximately known or drifters are
588 launched with a certain delay, using LAVA is still advantageous provided that
589 drifters are launched at a relative distance between each other of order R_d . In
590 this way, the target of interest in case of SAR or pollutant spill is expected to be
591 at a maximum distance of the order of half R_d from the drifters used in LAVA,
592 and the LAVA correction is expected to be significant.

593 *4.4. Discussion on forecasting skills*

594 All the results showed so far concern hindcast or nowcast applications. In
595 other words, if with $[t_0, t_1]$ we indicate the time period over which model, radar
596 and Lagrangian data are available, reconstructed LAVA trajectories and velocity
597 fields can be provided in the interval $[t_0, t_1]$ but not for times larger than t_1 .
598 Here we test whether or not the optimized LAVA fields at time t_0 have some
599 forecasting skills that can be used at least as a zero order approximation in
600 operational situations. To test this question, the original and LAVA blended
601 velocity fields at t_0 are held constant over the following 24 h and trajectories
602 are computed using this frozen field and compared with observed drifters.

603 Results are shown in Fig.8 for radar (left panel) and model (right panel) in
604 terms of $D(t)$ and $D_0(t)$. The pink (red) lines for radar (model) are obtained
605 using the original fields while the green (blue) lines are from setups similar
606 to E_R (E_M) experiments but with constant velocities. The radar and LAVA
607 blended results grow with the same trend and they reach 2-4 km uncertainty on
608 particle position during the first 6-10 h, even though LAVA results show slightly
609 reduced $D(t)$. Nevertheless, both curves are 1 to 2 km smaller than $D_0(t)$
610 for the whole 24 h period, indicating that the radar and LAVA blended fields
611 contribute to improve Lagrangian estimates. For the model, instead, LAVA
612 leads to a significant reduction on particle position uncertainty with respect to
613 the unblended model results. Within the very first 4 h the difference between

614 the two curves is about 2 km and it grows up to 8 km during the whole 24 h
615 period. This result reflects the effectiveness of the LAVA blending on the model,
616 as previously evidenced also by the Eulerian analysis. Overall, both radar and
617 model results show values of $D(t)$ of ≈ 2 km during the first 6 hours while $D_0(t)$
618 quickly reaches approximatively double values of 4 km. Therefore this very
619 simple forecast approximation is advantageous during the first 6-10 h for both
620 radar and model with respect to the zero order information represented by D_0 .
621 For the model in particular, using the LAVA blended field is especially useful
622 and allows to improve the results with respect to the original fields.

623 It should be noted that these results are expected to be dependent on the
624 correlation time scales of the Eulerian velocity field. A boundary current like
625 the one considered here might have longer correlation times and therefore higher
626 forecasting skills than for instance local coastal flows in gulfs or flows with large
627 tidal fluctuations. In this case, tidal currents are extremely low (generally lower
628 than 10^{-3} m/s) with amplitudes among the smallest in the whole Mediterranean
629 Sea [Alb erola et al., 1995b; Arabelos et al., 2011]. Even though in this paper tidal
630 effects have been disregarded for the aforementioned reasons, we also envision
631 possible applications of LAVA to marine basins where high frequency current
632 variability is large. This point will be further discussed in Section 5.

633 **5. Summary and concluding remarks**

634 In this paper, an extensive study on estimates of Lagrangian transport based
635 on radar, model and drifter data is presented. The study was performed during
636 the TOSCA experiment in the Toulon region. A number of diagnostics are used,
637 some of them of Eulerian nature, but the most relevant one is the Lagrangian
638 diagnostic $D(t)$ that computes the distance between drifter trajectories, consid-
639 ered as proxy for substances advected by the currents, and numerical trajec-
640 tories computed from velocity fields. The $D(t)$ diagnostic provides a quantitative
641 measure of the uncertainty on Lagrangian transport, and it is computed for
642 the original radar and model fields and for the LAVA blended ones. $D(t)$ is

643 compared with the measure of drifter absolute dispersion $D_0(t)$, that quantifies
644 how far drifters have traveled during the time t and that can be considered as
645 the uncertainty corresponding to the case of no available information on the
646 velocity fields.

647 Results show that the original radar fields reproduce well the mesoscale
648 pattern in the area and are therefore able to provide satisfactory estimates of
649 Lagrangian transport. $D(t)$ reaches approximately 6 km after 24 h while $D_0(t)$
650 is ≈ 12 km, indicating that the uncertainty is approximately halved using radar
651 velocities. The situation is very different for model results, that are character-
652 ized by $D(t)$ of the same order as $D_0(t)$ indicating a non significant reduction
653 of the uncertainty. When radar and model fields are blended with drifter data
654 through LAVA, in both cases $D(t)$ decreases significantly, indicating that the
655 uncertainty is strongly reduced. In quantitative terms, for LAVA experiments
656 with 7 blended drifters, $D(t)$ is of the order of 2 km after 24 h for both radar
657 and model. We notice that model results shown here might be a “worst case
658 scenario”, due for instance to a time lag error of the model in describing a prop-
659 agating feature as it can easily occur in highly non-linear flows in absence of
660 assimilation. The important point is that even in this difficult situation, trans-
661 port estimates can be greatly improved using LAVA. For the HF radar, the
662 improvement in the blended results is likely to be due to the fact that drifter
663 information allow to restore part of the smoothed variability, sharpening the
664 horizontal shear of the coastal jet. HF radar and drifters are highly correlated
665 and they provide a better estimate of transport compared to the model without
666 assimilating observations.

667 A sensitivity study is also performed considering the mean distance $\overline{d_L}$ be-
668 tween the drifters used in the LAVA blending and the control drifters used to
669 compute $D(t)$, and comparing it with the analysis scale $R \approx 7$ km which is in the
670 order of the typical Rossby radius R_d in the area. For the main configuration,
671 $\overline{d_L}$ is ≈ 250 m, i.e. $\overline{d_L} \ll R$, indicating that Lagrangian transport is computed
672 for particles initially close to the ones used for blending. Two sensitivity tests
673 are performed with $\overline{d_L} \approx R/4$ and $R/2$, respectively. $D(t)$ increases at increas-

674 ing $\overline{d_L}$ as expected, but maximum values stay below 4 km after 24 h, indicating
675 a persistent significant advantage in using LAVA. This is especially true for the
676 model where the uncertainty is still less than a third than for the original fields.

677 Finally, an investigation is carried out regarding possible forecasting skills.
678 The velocity fields are assumed constant in time starting from a certain time
679 t_0 and trajectories for the following period are computed using these frozen
680 fields and compared to drifter trajectories. The results show that, especially for
681 the model, there is a clear advantage in using LAVA blended fields, with the
682 uncertainty that is approximately half D_0 for a period of ≈ 6 h.

683 Overall, results show that estimates of Lagrangian transport are signifi-
684 cantly enhanced using LAVA blended fields. These results have a number of
685 consequences in terms of general understanding of Lagrangian predictability in
686 coastal flows, while also providing practical indications that can be useful for
687 operational purposes.

688 From the predictability point of view, results provide quantitative informa-
689 tion on the relevant space L and time T scales, that complement previous results
690 obtained in other regions and from numerical models [Griffa et al., 2004; Tail-
691 landier et al., 2006a]. Regarding the space scale L , the results on $\overline{d_L}$ sensitivity
692 indicate that the Rossby radius R_d is indeed the main parameter for Lagrangian
693 predictability. This is not obvious a priori, since R_d is a typical Eulerian scale
694 that characterizes the size of the most prominent features in the flow. Given
695 the high sensitivity of Lagrangian transport to the details of the flow, it could
696 be thought in principle that Lagrangian predictability scales L could be smaller
697 than the Rossby radius. Our results indicate instead that R_d is indeed relevant.

698 Regarding the time scale T , results on the forecasting skills indicate that
699 Lagrangian predictability is characterized by times of the order of a few hours
700 ($\approx 6-10$), typically a fraction of a day. This is much smaller than the Eulerian
701 time scale T_E that characterizes the persistency of the features and that in our
702 region is of the order of 3-6 days [Sammari et al., 1995]. T is likely to be related
703 to the typical Lagrangian time scale T_L that characterizes Lagrangian velocity
704 autocorrelation [Bauer et al., 2002]. The relationship between T_L and T_E has

705 been studied in a number of previous papers [Lumpkin et al., 2002; Middleton,
706 1985; Salle et al., 2008], and is generally quite complex and dependent on the
707 characteristics of the flow. When the flow is dominated by mesoscale structures,
708 in the so-called “frozen-field approximation”, T_L is dominated by advection
709 processes, and typically $T_L \ll T_E$. When instead the flow is dominated by
710 wind forcing or by tidal or inertial fluctuations, then $T_L \approx T_E$ and the so-called
711 “fixed float approximation” holds. The flow field considered here is likely to
712 be in-between the two approximations, since it is characterized by a boundary
713 current with strong mesoscale variability [Sammari et al., 1995] but it is also
714 subject to wind variations [Millot and Wald, 1980; Piterbarg et al., 2014] and
715 inertial oscillations [Petrenko, 2003].

716 From the practical and operational point of view, the results have a num-
717 ber of interesting implications. In the case of SAR or pollutant detection, the
718 improvement in transport estimates using LAVA could considerably change the
719 range of the search. While in case of no information the range is expected to be
720 of the order of D_0 , in case of LAVA blended fields the range decreases, for in-
721 stance to approximately $D_0/6$ in our experiments with 7 drifters for both radar
722 and models. Also, if radar and model velocities are not available in the area,
723 drifter data alone can be used to directly reconstruct the velocity field with
724 satisfactory results, at least where the coverage is appropriate.

725 The results also provide suggestions on drifter sampling. When the accident
726 location is known with accuracy, the best practice appears to be launching the
727 drifters as close as possible to the location in space and time. When instead the
728 accident location is not well known, as for instance for many SAR cases, then the
729 best launching practice is to cover the region of interest with a grid size of the
730 order of R_d . Even in this case, the LAVA blending can reduce the uncertainty
731 to approximately $D_0/3$. Also, the method reveals some forecasting skills, so
732 that LAVA blended fields can be used for operational purposes in a time range
733 of approximately 6 h, still providing an uncertainty decrease of approximately
734 $D_0/2$.

735 One important point is how these results can be generalized, especially re-

736 garding forecasting skills. As discussed earlier, the Lagrangian predictability
737 scale T is expected to be related to the flow characteristics in terms of T_L and
738 T_E . LAVA corrections and forecasting skills are expected to be more efficient
739 in flows with a persistent mesoscale component. In cases when the flow has a
740 large high frequency component, it might be useful to separate the two com-
741 ponents through filtering and applying LAVA only to the mesoscale part. This
742 is especially true when the high frequency component can be described deter-
743 ministically as for instance for a tidal flow. In this case, LAVA can be used
744 to increase mesoscale Lagrangian predictability while the high frequency part
745 can be deterministically superimposed [Taillandier et al., 2006a] even though er-
746 rors in the tidal component can occur and associated dispersion might be only
747 partially captured.

748 A number of directions for future work can be envisioned. Here we have
749 focused on the LAVA blending technique, that can be easily used in operational
750 settings and takes advantage of any velocity field available, but that has re-
751 stricted forecasting skills. Full assimilation of surface velocity data from HF
752 radars or drifters, on the other hand, is expected to provide more extended
753 forecasts, even though from the operational point of view it is more restricted
754 since the model has to be appropriately set up in advance for the area. In the
755 future, it would be useful to optimize the use of both approaches in operational
756 settings. Also, specific technical approaches in the assimilation of surface data
757 should be compared and further investigated. As an example, some efforts so far
758 have concentrated on using filtered data that represent low frequency dynamics
759 [Lipphardt et al., 2000; Oke et al., 2002], while other approaches have focused
760 on identifying the main parameters to be optimized in a given area, such as
761 wind forcing or boundary conditions [Barth et al., 2010, 2011; Marmain et al.,
762 2014; Paduan and Shulman, 2004; Shulman and Paduan, 2009]. The LAVA
763 technique could also be used to reduce radar uncertainty in emergency cases or
764 during a rapid response operation, when radar antennas cannot be previously
765 calibrated due to lack of time. An other important direction for further investi-
766 gations is the study of optimal sampling design for drifter launchings. Here we

767 have mostly investigated the impact on LAVA results, but it should be noted
768 that in many practical applications other information for instance in terms of
769 relative dispersion is desirable from closely launched clusters [Schroeder et al.,
770 2012; Haza et al., 2013]. Optimized launchings should then combine information
771 from various scales and different metrics. And finally, a very important issue for
772 future studies is the improvement of drifter designs in order to make them more
773 eco-friendly and biodegradable. This could really open new scenarios, allowing
774 for the use of extensive drifter deployments also for operational purposes, such
775 as a GPS-tracked version of the compact drift-card type instrument.

776 *Acknowledgements*

777 The authors gratefully acknowledge support from the MED TOSCA project, co-
778 financed by the European Regional Development Fund. The GLAZUR64 simulations
779 were performed using NEC-SX8 from GENCI-IDRIS resources (Grant 2012011707).
780 The authors also wish to thank Mercator Ocean and Météo-France for providing nu-
781 merical data. The MIO RADAR team (Y. Barbin, J. Gaggelli, C. Quentin, P. Guter-
782 man) is also acknowledged for its major contribution to radar field campaigns of this
783 study. M. Berta, M. G. Magaldi, A. Griffa and M. Borghini also acknowledge support
784 from the Italian Flagship Project RITMARE. The authors also thank two anonymous
785 reviewers for improving the manuscript.

786 **References**

787 Albérola, C., Millot, C., Font, J., 1995a. On the seasonal and mesoscale variabil-
788 ities of the Northern Current during the PRIMO-0 experiment in the western
789 Mediterranean Sea. *Oceanologica Acta* 18, 163–192.

- 790 Albérola, C., Rousseau, S., Millot, C., Astraldi, M., Font, J., Garcia-Lafuente,
791 J., Gasparini, G.P., Send, U., Vangriesheim, A., 1995b. Tidal currents West-
792 ern Mediterranean. *Oceanol. Acta* 18, 273–284.
- 793 Arabelos, D.N., Papazachariou, D.Z., Contadakis, M.E., Spatalas, S.D., 2011. A
794 new tide model for the Mediterranean Sea based on altimetry and tide gauge
795 assimilation. *Oc. Sci.* 7, 429–444.
- 796 Barth, A., Alvera-Azcárate, A., Beckers, J.M., Staneva, J., Stanev, E.V.,
797 Schulz-Stellenfleth, J., 2011. Correcting surface winds by assimilating High-
798 Frequency radar surface currents in the German Bight. *Ocean Dyn.* 61, 599–
799 610.
- 800 Barth, A., Alvera-Azcárate, A., Gurgel, K.W., Staneva, J., Port, A., Beckers,
801 J.M., Stanev, E.V., 2010. Ensemble perturbation smoother for optimizing
802 tidal boundary conditions by assimilation of High-Frequency radar surface
803 currents-application to the German Bight. *Ocean Science* 6.
- 804 Bauer, S., Swenson, M.S., Griffa, A., 2002. Eddy mean flow decomposition
805 and eddy diffusivity estimates in the tropical Pacific Ocean: 2. results. *J.*
806 *Geophys. Res.* 107, 3154–3171.
- 807 Bellomo, L., Berta, M., Gasparini, G., Griffa, A., Magaldi, M., Marmain, J.,
808 Molcard, A., Vetrano, A., Béguey, L., Borghini, M., 2013. Observational
809 evidence of mesoscale variability of the Northern Current (North-Western
810 Mediterranean Sea): a combined study via gliders, HF RADAR, surface
811 drifters, and vessel data, in: EGU General Assembly, Vienna (Austria). pp.
812 EGU2013–5469–2.
- 813 Bosse, A., Testor, P., Mortier, L., Béguey, L., Bernardet, K., Taillandier, V.,
814 dOrtenzio, F., Prieur, L., Coppola, L., Bourrin, F., 2013. New insights of the
815 Northern Current in the Western Mediterranean Sea from Gliders data: Mean
816 structure, Transport, and Seasonal Variability, in: EGU General Assembly,
817 Vienna (Austria). pp. EGU2013–11315–2.

- 818 Chang, Y., Hammond, D., Haza, A., Hogan, P., Huntley, H., Jr., A.K., Jr.,
819 B.L., Taillandier, V., Griffa, A., Ozgokmen, T., 2011. Enhanced estimation
820 of sonobuoy trajectories by velocity reconstruction with near-surface drifters.
821 *Oc. Model.* 36, 179 – 197.
- 822 Chapman, R., Graber, H., 1997. Validation of hf radar measurements. *Oceanog-*
823 *raphy* 10, 76–79.
- 824 Chapman, R.D., Shay, L.K., Graber, H.C., Edson, J.B., Karachintsev, A.,
825 Trump, C.L., Ross, D.B., 1997. On the accuracy of HF radar surface current
826 measurements: Intercomparisons with ship-based sensors. *J. Geophys. Res.*
827 102, 737–748.
- 828 Davis, R.E., 1985. Drifter observations of coastal surface currents during CODE:
829 The method and descriptive view. *J. Geophys. Res.* 90, 4741–4755.
- 830 Derber, J., Rosati, A., 1989. A global oceanic data assimilation system. *J. Phys.*
831 *Oceanogr.* 19, 1333–1347.
- 832 Edwards, K., Werner, F., Blanton, B., 2006. Comparison of observed and mod-
833 eled drifter trajectories in coastal regions: An improvement through adjust-
834 ments for observed drifter slip and errors in wind fields. *J. Atmos. Ocean*
835 *Tech.* 23, 1614–1620.
- 836 Emery, B.M., Washburn, L., Harlan, J.A., 2004. Evaluating Radial Current
837 Measurements from CODAR High-Frequency Radars with Moored Current
838 Meters. *J. Atmos. Ocean. Technol.* 21, 1259–1271.
- 839 Essen, H.H., Gurgel, K.W., Schlick, T., 2000. On the accuracy of current mea-
840 surements by means of HF radar. *IEEE J. Oceanic Eng.* 25, 472–480.
- 841 Griffa, A., Piterbarg, L.I., Ozgokmen, T., 2004. Predictability of Lagrangian
842 particle trajectories: Effects of smoothing of the underlying Eulerian flow. *J.*
843 *Marine Res.* 62, 1–35.

- 844 Guihou, K., , Marmain, J., Ourmières, Y., Molcard, A., Zakardjian, B., Forget,
845 P., 2013. A case study of the mesoscale dynamics in the North-Western
846 Mediterranean Sea: a combined datamodel approach. *Ocean Dyn.* 63, 793–
847 808.
- 848 Gurgel, K.W., Antonischki, G., Essen, H.H., Schlick, T., 1999a. Wellen Radar
849 (WERA): a new ground-wave HF radar for ocean remote sensing. *Coastal*
850 *Engineering* 37, 219–234.
- 851 Gurgel, K.W., Barbin, Y., 2008. Suppressing radio frequency interference in HF
852 radars. *Sea Technology* 49, 39–42.
- 853 Gurgel, K.W., Essen, H.H., Kingsley, S.P., 1999b. High-Frequency radars: Phys-
854 ical limitations and recent developments. *Coastal Engineering* 37, 201–218.
- 855 Haller, G., Poje, A., 1998. Finite time transport in aperiodic flow. *Physica D*
856 119, 352–380.
- 857 Harlan, J., Terrill, E., Hazard, L., Keen, C., Barrick, D., Whelan, C., Howden,
858 S., Kohut, J., 2010. The integrated ocean observing system high-frequency
859 radar network: status and local, regional and national applications. *Mar.*
860 *Technol. Soc. J.* 44, 122–132.
- 861 Haza, A., Poje, A., Ozgokmen, T., Griffa, A., Haus, B., Huntley, H., Hogan, P.,
862 Jacobs, G., Kirwan Jr., A., Lipphardt, B., Novelli, G., Olascoaga, J., Beron-
863 Vera, F., Reniers, A., Ryan, E., 2013. Surface relative dispersion measure-
864 ments in the CARTHE’s grand lagrangian deployment (GLAD) experiment.,
865 in: 2013 Gulf of Mexico Oil Spill & Ecosystem Science Conference - New
866 Orleans (USA).
- 867 Haza, A.C., Griffa, A., Martin, P., Molcard, A., Ozgokmen, T.M., Poje, A.C.,
868 Barbanti, R., Book, J.W., Poulain, P.M., Rixen, M., Zanasca, P., 2007.
869 Model-based directed drifter launches in the Adriatic Sea: Results from the
870 DART experiment. *Geophys. Res. Lett.* 34.

- 871 Kaplan, D.K., Largier, J., Botsford, L.W., 2005. HF radar observations of
872 surface circulation off Bodega Bay (northern California, USA). *J. Geophys.*
873 *Res.* 110, C10020.
- 874 Kohut, J., Glenn, S., 2003. Improving HF radar surface current measurements
875 with measured antenna beam patterns. *J. Atmos. Ocean Tech.* 20, 1303–1316.
- 876 Kohut, J., Roarty, H., Randall-Goodwin, E., Glenn, S., Lichtenwalner, C., 2012.
877 Evaluation of two algorithms for a network of coastal HF radars in the Mid-
878 Atlantic Bight. *Oc. Dyn.* 62, 953–968.
- 879 Krause, P., Restrepo, J., 2009. The diffusion kernel filter applied to Lagrangian
880 data assimilation. *Mon. Wea. Rev.* 137, 4386–4400.
- 881 Kuang, L., Blumberg, A., Georgas, N., 2012. Assessing the fidelity of surface
882 currents from a coastal ocean model and HF radar using drifting buoys in the
883 Middle Atlantic Bight. *Ocean Dynam.* 62, 1229–1243.
- 884 Kuznetsov, L., Ide, K., Jones, C., 2009. A method for assimilation of Lagrangian
885 data. *Mon. Wea. Rev.* 131, 2247–2260.
- 886 Lipa, B., Nyden, B., Ullman, D.S., Terrill, E., 2006. SeaSonde Radial Velocities:
887 Derivation and Internal Consistency. *IEEE Journal of Oceanic Engineering*
888 31, 850–861.
- 889 Lipa, B.J., Barrick, D.E., 1983. Least-Squares Methods for the Extraction of
890 Surface Currents from CODAR Crossed-Loop Data: Application at ARSLOE.
891 *IEEE Journal of Oceanic Engineering* OE-8, 226–253.
- 892 Lipphardt, B.L., Kirwan, A.D., Grosch, C.E., Lewis, J.K., Paduan, J.D., 2000.
893 Blending HF radar and model velocities in Monterey Bay through normal
894 mode analysis. *J. Geophys. Res.* 105, 3425–3450.
- 895 Lumpkin, R., Treguier, A., Speer, K., 2002. Lagrangian eddy scales in the
896 Northern Atlantic Ocean. *J. Phys. Oceanogr.* 32, 2425–2440.

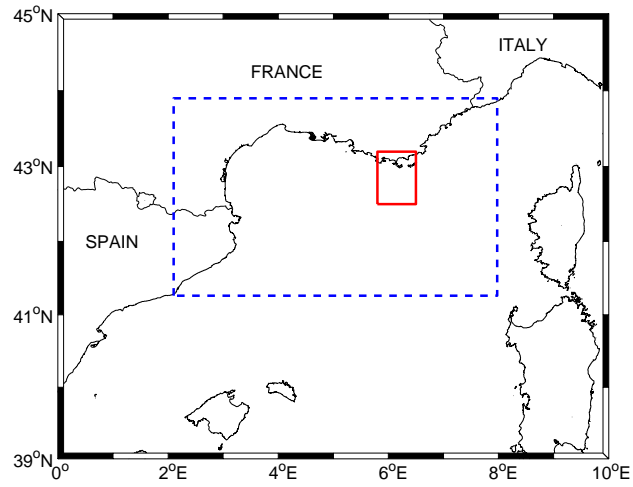
- 897 Madec, G., 2008. NEMO ocean engine, in: Technical Report, Institut Pierre-
898 Simon Laplace, France. p. Note du Pole de modelisation.
- 899 Mantovanelli, A., Heron, M., Prytz, A., Steinberg, C., Wisdom, D., 2011. Vali-
900 dation of radar-based lagrangian trajectories against surface-drogued drifters
901 in the Coral Sea, Australia, in: OCEANS 2011, pp. 1–4.
- 902 Marmain, J., Molcard, A., Forget, P., Barth, A., 2014. Assimilation of HF radar
903 surface currents to optimize forcing in the North Western Mediterranean Sea.
904 *Nonlinear Proc. Geoph.* 21, 659–675.
- 905 Marullo, S., Salusti, E., Viola, A., 1985. Observation of a small-scale baroclinic
906 eddy in the Ligurian Sea. *Deep-Sea Res.* 32, 215–222.
- 907 Middleton, J.F., 1985. Drifter spectra and diffusivities. *J. Marine Res.* 43,
908 37–55.
- 909 Millot, C., Wald, L., 1980. The effect of Mistral wind on the Ligurian current
910 near Provence. *Oceanol. Acta* 3, 399–402.
- 911 Molcard, A., Piterbarg, L.I., Griffa, A., Ozgokmen, T.M., Mariano, A.J., 2003.
912 Assimilation of drifter observations for the reconstruction of the Eulerian
913 circulation field. *J. Geophys. Res.* 108.
- 914 Molcard, A., Poulain, P., Forget, P., Griffa, A., Barbin, Y., Gaggelli, J., Maistre,
915 J.D., Rixen, M., 2009. Comparison between vhf radar observations and data
916 from drifter clusters in the gulf of la spezia (mediterranean sea). *Journal of*
917 *Marine Systems* 78, S79–S89.
- 918 Ohlmann, C., White, P., Washburn, L., Emery, B., Terrill, E., Otero, M.,
919 2007. Interpretation of coastal HF radar-derived surface currents with high-
920 resolution drifter data. *J. Atmos. Ocean Tech.* 24, 666–680.
- 921 Oke, P.R., Allen, J.S., Miller, R.N., Egbert, G.D., Kosro, P.M., 2002. Assimi-
922 lation of surface velocity data into a primitive equation coastal ocean model.
923 *J. Geophys. Res.* 107.

- 924 Olascoaga, M., Rypina, I., Brown, M., Beron-Vera, F., Kocak, H., Brand, L.,
925 Halliwell, G., Shay, L., 2006. Persistent transport barrier on the West Florida
926 Shelf. *Geophys. Res. Lett.* 33.
- 927 Olascoaga, M.J., Haller, G., 2012. Forecasting sudden changes in environmental
928 pollution patterns. *P. Natl. Acad. Sci. USA* 109, 4738–4743.
- 929 Ourmières, Y., Zakardjian, B., Béranger, K., Langlais, C., 2011. Assessment of a
930 NEMO-based downscaling experiment for the North-Western Mediterranean
931 region: Impacts on the Northern Current and comparison with ADCP data
932 and altimetry products. *Ocean Mod.* 39, 386–404.
- 933 Paduan, J., Kim, K.C., Cook, M., Chavez, F., 2006. Calibration and validation
934 of direction-finding high-frequency radar ocean surface current observations.
935 *IEEE J. Oceanic Eng.* 31, 862–875.
- 936 Paduan, J.D., Rosenfeld, L.K., 1996. Remotely sensed surface currents in Mon-
937 terey Bay from shore based HF radar (Coastal Ocean Dynamics Application
938 Radar). *J. Geophys. Res.* 101, 20669–20686.
- 939 Paduan, J.D., Shulman, I., 2004. HF radar data assimilation in the Monterey
940 Bay area. *J. Geophys. Res.* 109.
- 941 Paduan, J.D., Washburn, L., 2013. High-frequency radar observations of ocean
942 surface currents. *Annu. Rev. Mar. Sci.* 5, 115–136.
- 943 Petrenko, A.A., 2003. Variability of circulation features in the Gulf of Lion
944 NW Mediterranean Sea. Importance of inertial currents. *Oceanol. Acta* 26,
945 323–338.
- 946 Picco, P., Cappelletti, A., Sparnocchia, S., Schiano, M., Pensieri, S., Bozzano,
947 R., 2010. Upper layer current variability in the Central Ligurian Sea. *Ocean*
948 *Sci.* 6, 825–836.
- 949 Piterbarg, L., Taillandier, V., Griffa, A., 2014. Investigating frontal variability
950 from repeated glider transects in the Ligurian Current (North West Mediter-
951 ranean Sea). *J. Marine Syst.* 129, 381–395.

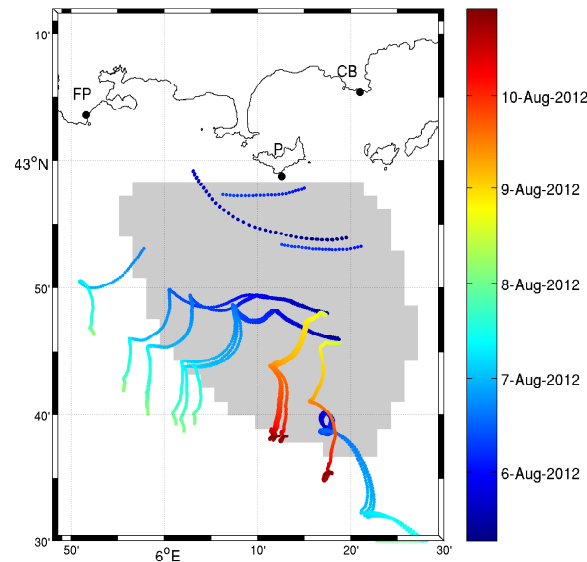
- 952 Poulain, P.M., 1999. Drifter observations of surface circulation in the Adriatic
953 Sea between December 1994 and March 1996. *J. Marine Syst.* 20, 231 – 253.
- 954 Poulain, P.M., Gerin, R., Mauri, E., Pennel, R., 2009. Wind effects on drogued
955 and undrogued drifters in the Eastern Mediterranean. *J. Atmos. Ocean Tech.*
956 26, 1144–1156.
- 957 Poulain, P.M., Ursella, L., Brunetti, F., 2002. Direct measurements of water-
958 following characteristics of CODE surface drifters, in: LAPCOD Meeting.
- 959 Pugh, D.T., 1987. *Tides, Surges and Mean Sea-Level.* John Wiley and Sons,
960 New York.
- 961 Robinson, A., Leslie, W., Theocharis, A., Lascaratos, A., 2001. Mediterranean
962 Sea circulation, in: Steele, J.H., Thorpe, S.A., Turekian, K.K. (Eds.), *Ocean
963 currents: a derivative of the encyclopedia of ocean sciences.* Academic press
964 San Diego.
- 965 Rypina, I.I., Kirincich, A.R., Limeburner, R., Udovydchenkov, I.A., 2014. Eu-
966 lerian and Lagrangian Correspondence of High-Frequency Radar and Surface
967 Drifter Data: Effects of Radar Resolution and Flow Components. *J. Atmos.
968 Ocean. Technol.* 31, 945–966.
- 969 Salle, J.B., Speer, K., Morrow, R., Lumpkin, R., 2008. An estimate of La-
970 grangian eddy statistics and diffusion in the mixed layer of the Southern
971 Ocean. *J. Marine Res.* 66, 441–463.
- 972 Sammari, C., Millot, C., Prieur, L., 1995. Aspects of the seasonal and mesoscale
973 variabilities of the Northern Current in the western Mediterranean Sea in-
974 ferred from the PROLIG-2 and PROS-6 experiments. *Deep-Sea Res.* 42,
975 893–917.
- 976 Schaeffer, A., Molcard, A., Forget, P., Fraunie, P., Garreau, P., 2011. Generation
977 mechanisms for mesoscale eddies in the Gulf of Lions: radar observation and
978 modeling, in: *Ocean Dyn.*, pp. 1587–1609.

- 979 Schroeder, K., Chiggiato, J., Haza, A.C., Griffa, A., Ozgokmen, T.M., Zanasca,
980 P., Molcard, A., Borghini, M., Poulain, P.M., Gerin, R., Zambianchi, E.,
981 Falco, P., Trees, C., 2012. Targeted Lagrangian sampling of submesoscale
982 dispersion at a coastal frontal zone. *Geophys. Res. Lett.* 39.
- 983 Sentchev, A., Forget, P., Barbin, Y., Yaremchuk, M., 2013. Surface circulation in
984 the Iroise Sea (W. Brittany) from high resolution HF radar mapping. *Journal*
985 *of Marine Systems* 109-110, S153–S168.
- 986 Shadden, S.C., Lekien, F., Paduan, J.D., Chavez, F.P., Marsden, J.E., 2009.
987 The correlation between surface drifters and coherent structures based on
988 high-frequency radar data in Monterey Bay. *Deep-Sea Res. Pt II* 56, 161 –
989 172.
- 990 Shay, L.K., Cook, T.M., Hallock, Z., Haus, B.K., Graber, H.C., Martinez, J.,
991 2001. The Strength of the m_2 Tide at the Chesapeake Bay Mouth. *J. Phys.*
992 *Oceanogr.* 31, 427–449.
- 993 Shay, L.K., Lee, T.N., Williams, E.J., Graber, H., Rooth, C.G.H., 1998a. Effects
994 of low-frequency current variability on submesoscale near-inertial vortices. *J.*
995 *Geophys. Res.* 103, 18691–18714.
- 996 Shay, L.K., Lentz, S.J., Graber, H.C., Haus, B.K., 1998b. Current Structure
997 Variations Detected by High-Frequency Radar and Vector-Measuring Current
998 Meters. *J. Atmos. Ocean. Technol.* 15, 237–256.
- 999 Shay, L.K., Martinez-Pedraja, J., Cook, T.M., Haus, B.K., Weisberg, R.H.,
1000 2007. High Frequency Radar Mapping of Surface Currents Using WERA. *J.*
1001 *Atmos. Ocean. Technol.* 24, 485–503.
- 1002 Shulman, I., Paduan, J.D., 2009. Assimilation of HF radar-derived radials and
1003 total currents in the Monterey Bay area. *Deep Sea Res. Pt. II* 56, 149 – 160.
- 1004 Stewart, R.H., Joy, J.W., 1974. HF radio measurements of surface currents.
1005 *Deep-Sea Res.* 21, 1039–1049.

- 1006 Taillandier, V., Dobricic, S., Testor, P., Pinardi, N., Griffa, A., Mortier, L., Gas-
1007 parini, G.P., 2010. Integration of Argo trajectories in the Mediterranean Fore-
1008 casting System and impact on the regional analysis of the western Mediter-
1009 ranean circulation. *J. Geophys. Res.* 115.
- 1010 Taillandier, V., Griffa, A., Molcard, A., 2006a. A variational approach for the
1011 reconstruction of regional scale Eulerian velocity fields from Lagrangian data.
1012 *Oc. Model.* 13.
- 1013 Taillandier, V., Griffa, A., Poulain, P.M., Branger, K., 2006b. Assimilation of
1014 Argo float positions in the north western Mediterranean Sea and impact on
1015 ocean circulation simulations. *Geophys. Res. Lett.* 33.
- 1016 Taillandier, V., Griffa, A., Poulain, P.M., Signell, R., Chiggiato, J., Carniel,
1017 S., 2008. Variational analysis of drifter positions and model outputs for the
1018 reconstruction of surface currents in the central Adriatic during fall 2002. *J.*
1019 *Geophys. Res.* 113.
- 1020 Ullman, D.S., O'Donnell, J., Kohut, J., Fake, T., Allen, A., 2006. Trajectory
1021 prediction using HF radar surface currents: Monte Carlo simulations of pre-
1022 diction uncertainties. *J. Geophys. Res.* 111.
- 1023 Weaver, A., Courtier, P., 2001. Correlation modelling on the sphere using a
1024 generalized diffusion equation. *Quart. J. Roy. Meteor. Soc.* 127, 1815–1846.



(a)



(b)

Figure 1: (a) Western Mediterranean basin. The dashed blue line represents the limits of the GLAZUR64 model domain. The red solid square indicates the area of the experiment in front of Toulon (France), expanded in panel (b) where the gray area is the radar coverage, the black dots show the radar sites (FP-Fort Peyras, P-Porquerolles and CB-Cap Bénat). The whole set of 20 drifter trajectories are color coded for time during the period from 5-Aug-2012 06:41:00 to 10-Aug-2012 22:21:00 (UTC times). The domain width is chosen to give an overview on drifters and close trajectories may not be distinguishable.

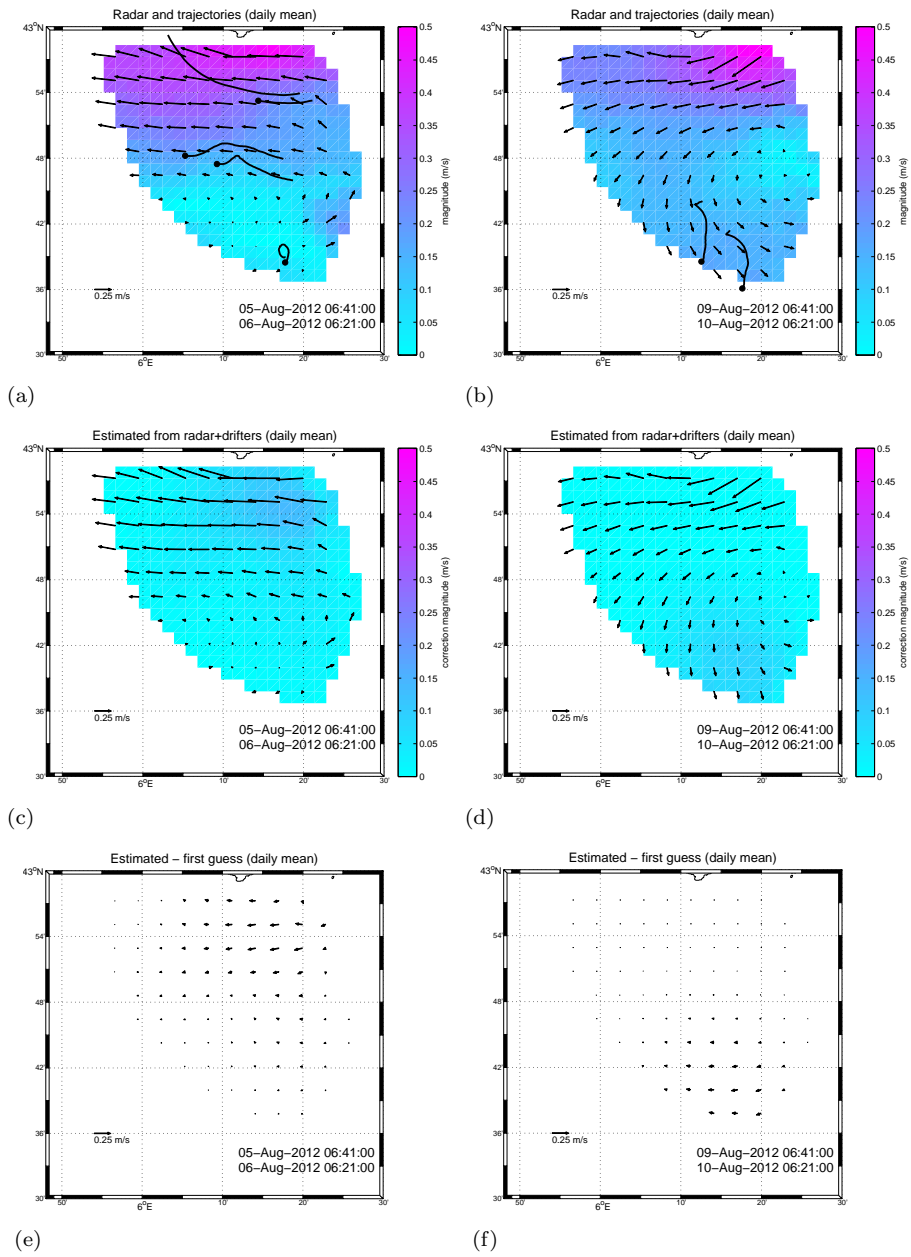


Figure 2: Examples of daily average radar velocities for August 5 (left panels) and August 9 (right panels). (a) and (b) show the original velocity fields, with color indicating velocity amplitude and arrows indicating velocity vectors. Drifter trajectories used in LAVA are superimposed; (c),(d) show the LAVA blended fields (E_R), with color indicating the amplitude of the correction C_{daily} (equation 1) and arrows indicating velocity vectors. Panels (e) and (f) represent the vectorial difference between the LAVA estimated and the original velocity field.

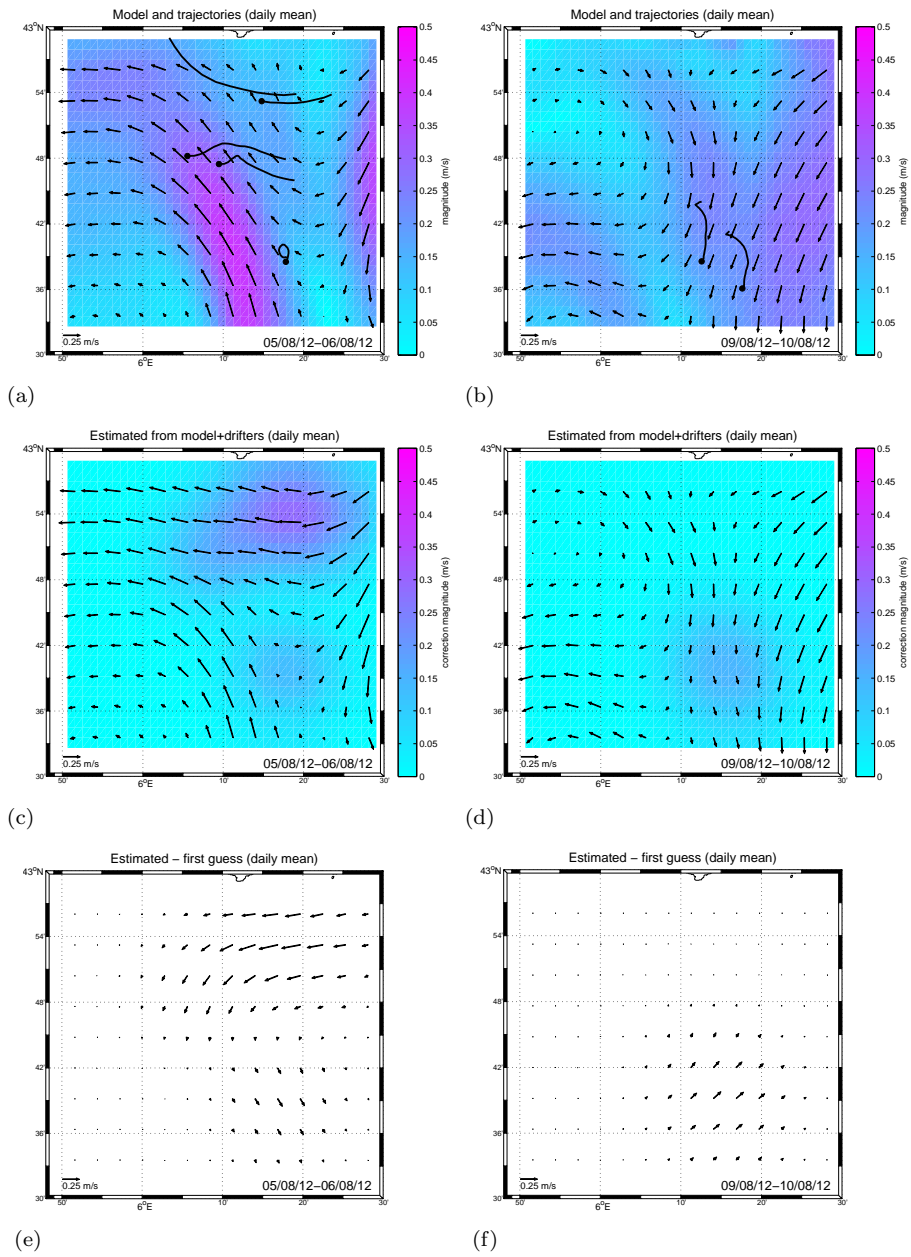


Figure 3: Examples of daily average model velocities for August 5 (left panels) and August 9 (right panels). (a) and (b) show the original velocity fields, with color indicating velocity amplitude and arrows indicating velocity vectors. Drifter trajectories used in LAVA are superimposed; (c),(d) show the LAVA blended fields (E_M), with color indicating the amplitude of the correction C_{daily} (equation 1) and arrows indicating velocity vectors. Panels (e) and (f) represent the vectorial difference between the LAVA estimated and the original velocity field.

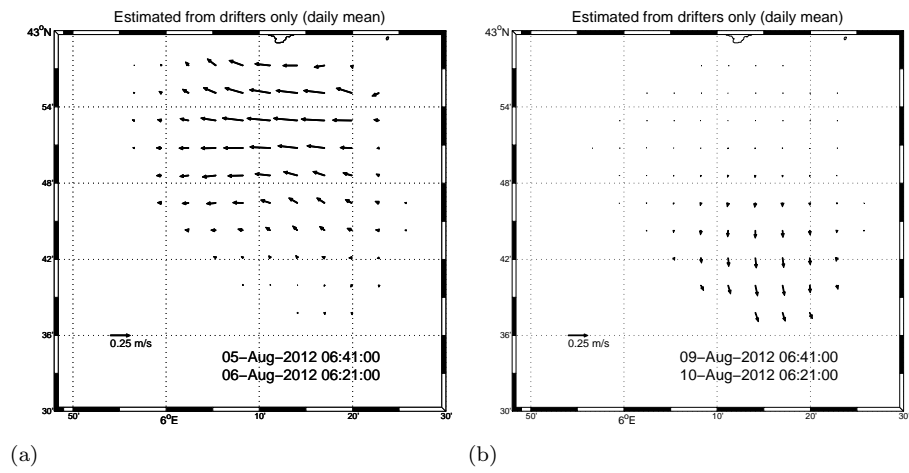


Figure 4: Examples of daily average velocities reconstructed from drifters only using LAVA (E_D) for August 5 (left panel) and August 9 (right panel). Arrows indicate velocity vectors.

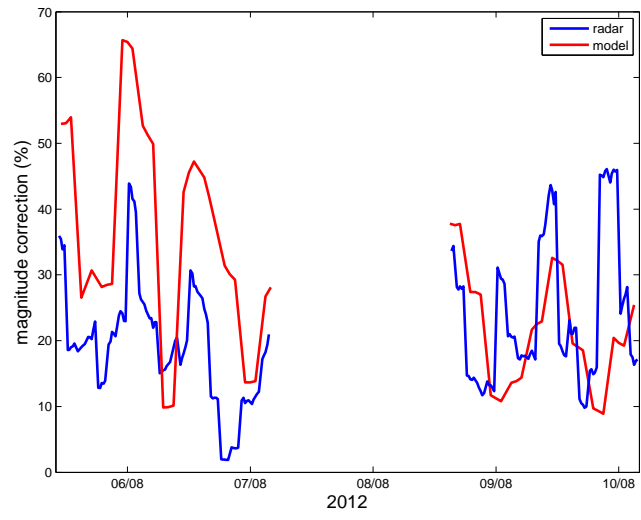


Figure 5: Time series of the normalized correction C_{norm} (equation 2) for radar (blue) and model (red).

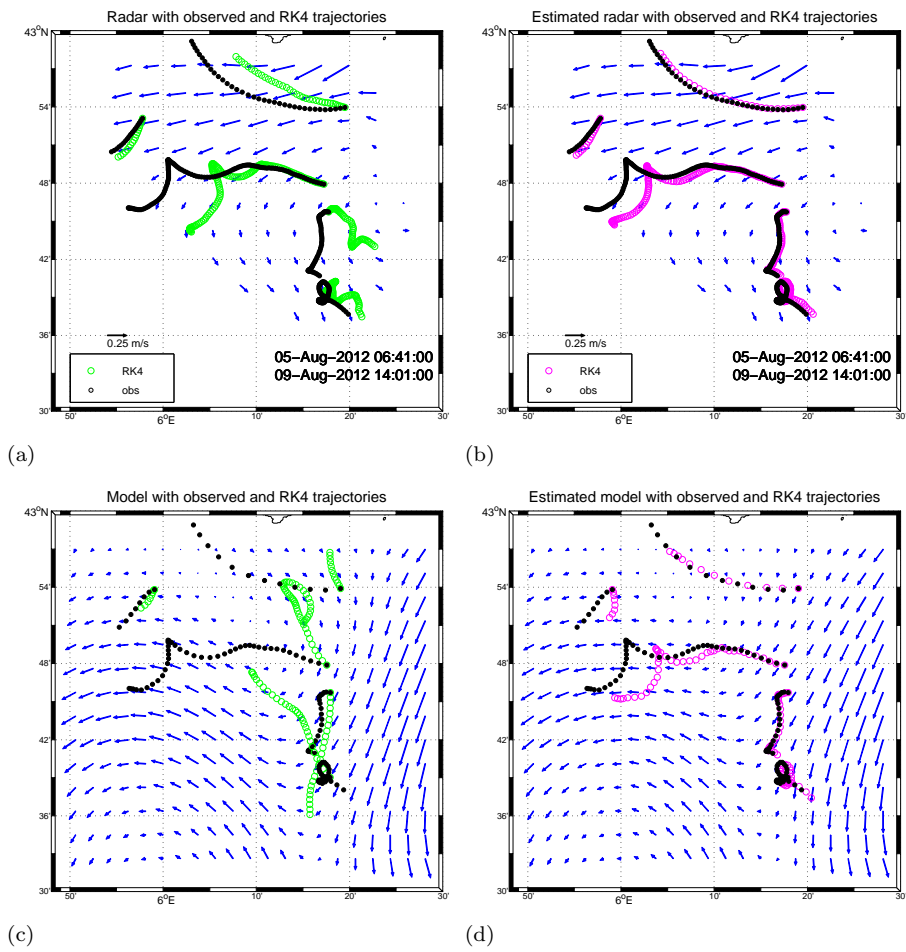


Figure 6: Comparison between observed drifter trajectories (black) and simulated trajectories computed from different velocity fields: (a) original radar velocity; (b) LAVA blended radar velocity; (c) original model velocity; (d) LAVA blended model velocity. The dots corresponds to the positions every Δt (see Table1). The trajectories are superimposed to the corresponding velocity averaged over the period 5-9 August. Arrows size and period considered in the lower panels are as in the upper ones.

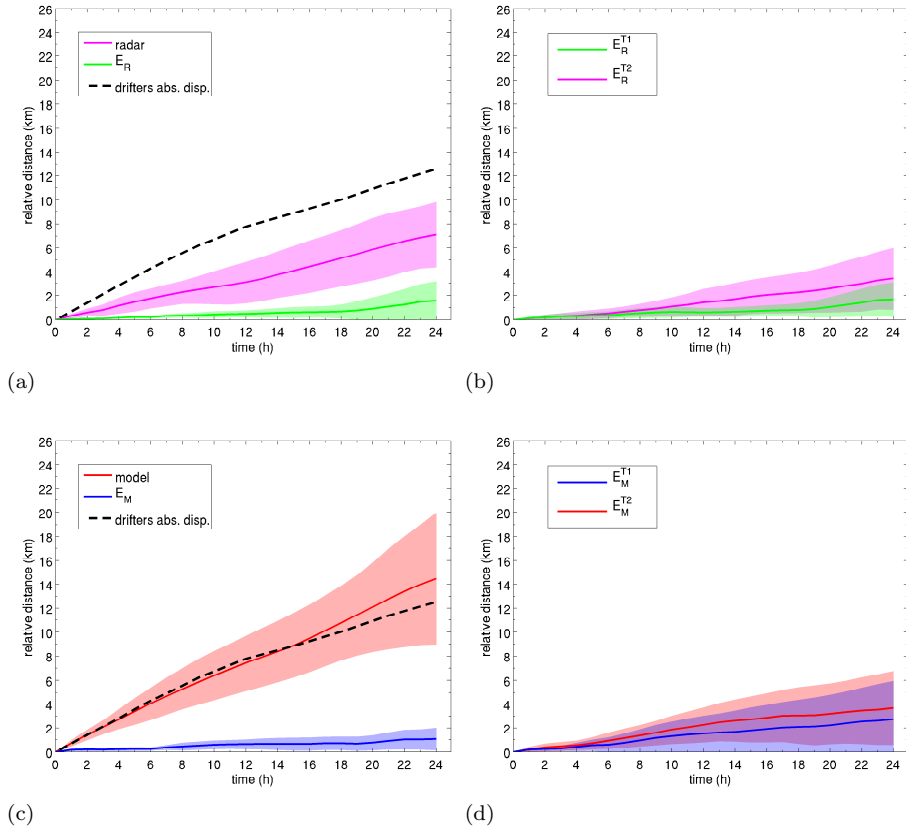


Figure 7: Time series of separation between observed and simulated trajectories $D(t)$ (equation 4) for various cases testing LAVA configurations. $D(t)$ is indicated by color lines (shades indicate standard deviations) and is compared to drifter absolute dispersion $D_0(t)$ (equation 5) indicated by black dashed lines. Upper (lower) panels are for radar (model) velocities. (a) pink line is computed for original radar velocity and green line is for the E_R experiment; (b) green and pink lines are computed for downgraded LAVA E_R^{T1} and E_R^{T2} experiments, respectively. (c) red line is computed for original model velocity and blue line for the E_M experiment; (d) blue and red lines are computed for downgraded LAVA E_M^{T1} and E_M^{T2} experiments, respectively.

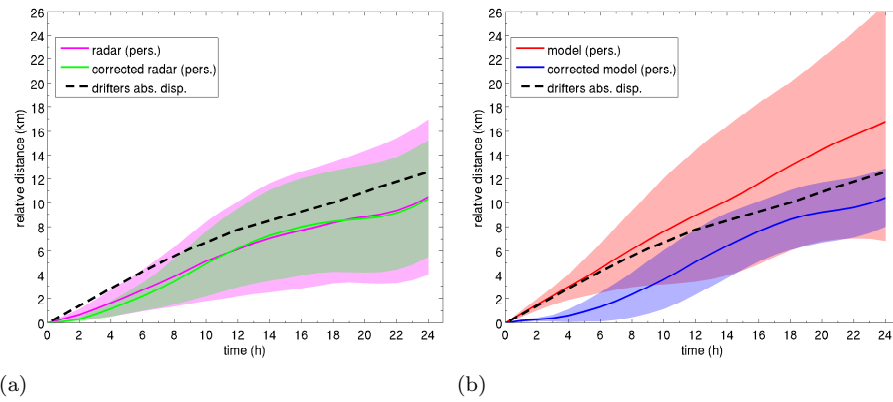


Figure 8: Time series of separation between observed and simulated trajectories $D(t)$ (equation 4) for various cases testing forecasting skills. The velocity field is held constant in time and trajectories are computed using the frozen field. $D(t)$ is indicated by color lines (shades indicate standard deviations) and is compared to drifter absolute dispersion $D_0(t)$ (equation 5) indicated by black dashed lines. Left (right) panel is for radar (model) velocities. (a) pink line is computed for original radar velocity and green line is for a setup like E_R but with constant radar velocities; (c) red line is computed for original model velocity and blue line is for a setup like E_M but with constant model velocities.

	Experiments						
	E_R	E_M	E_D	E_R^{T1}	E_R^{T2}	E_M^{T1}	E_M^{T2}
first guess velocity	Radar	Model	-	Radar	Radar	Model	Model
N_{Dft}	7	7	7	5	5	5	5
Δx	2 km	1/64°	2 km	2 km	2 km	1/64°	1/64°
Δt	20 min	1 h	20 min	20 min	20 min	1 h	1 h
R	7 km	7 km	7 km	7 km	7 km	7 km	7 km
T_a	2 h	4 h	2 h	2 h	2 h	4 h	4 h
d_L	247 m	256 m	247 m	1.16 km	3.41 km	1.49 km	3.34 km

Table 1: LAVA parameters used for all experiments (see Section 3.2).

	5-7 aug	8-10 aug
C_{norm} (radar, E_R) (%)	19.21	24.96
C_{norm} (model, E_M) (%)	36.17	21.56
drifters average velocity (m/s)	0.23	0.18

Table 2: Average values of C_{norm} (equation 2) computed during the two periods of drifter coverage (Fig.5) and compared to drifter velocity.